

# Dark Plasma and Cluster Mergers

Matti Heikinheimo

KBFI, Tallinn

[1504.04371 [hep-ph]], M. H., M. Raidal, C. Spethmann, H. Veermäe

21.05.2015

# Contents

## 1 Introduction

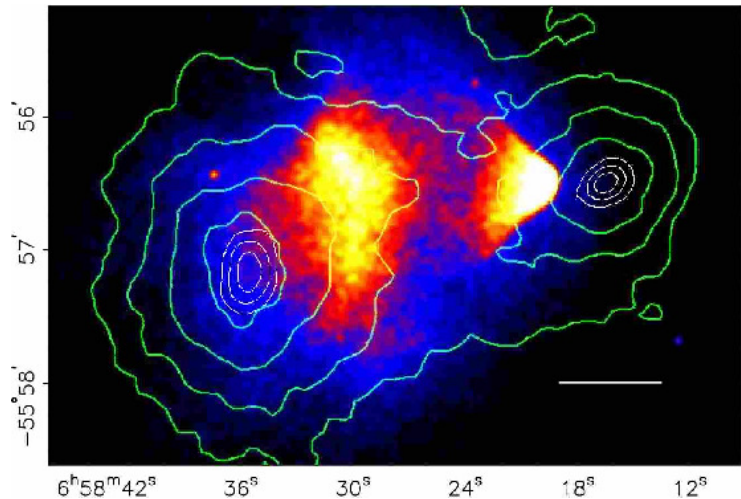
## 2 Plasma dynamics

## 3 The model

# Astrophysical observations

- **1E 0657-558:** in the bullet cluster, the dark matter halo of the subcluster is observed to pass through the main cluster [astro-ph/0309303].

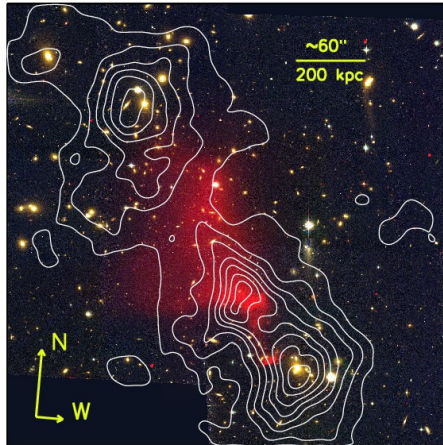
# 1E 0657-558



# Astrophysical observations

- **1E 0657-558:** in the bullet cluster, the dark matter halo of the subcluster is observed to pass through the main cluster [astro-ph/0309303].
- **Abell 520:** an excess of dark matter observed on top of the visible X-ray emitting gas, between the merging clusters [1401.3356 [astro-ph.CO]].

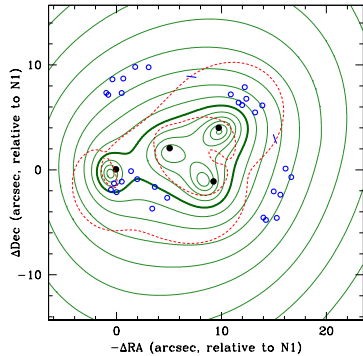
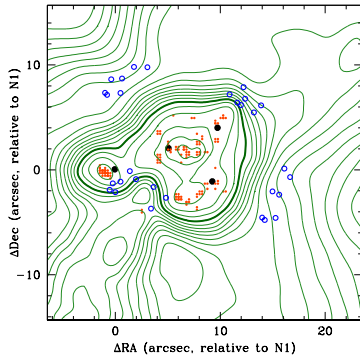
# Abell 520



# Astrophysical observations

- **1E 0657-558:** in the bullet cluster, the dark matter halo of the subcluster is observed to pass through the main cluster [astro-ph/0309303].
- **Abell 520:** an excess of dark matter observed on top of the visible X-ray emitting gas, between the merging clusters [1401.3356 [astro-ph.CO]].
- **Abell 3827:** a separation between the dark matter halo and the visible stars observed in the central four galaxies [1504.03388 [astro-ph.CO]].

# Abell 3827





# Contents

1 Introduction

2 Plasma dynamics

3 The model

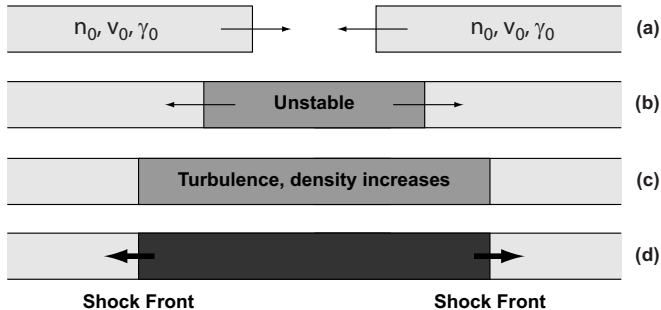
# Plasma

Plasma is a fluid, where

- the size of the fluid is large compared to the Debye shielding length  $\lambda_D = \sqrt{\frac{T}{4\pi\alpha n}}$  (bulk interactions dominate over surface effects),
- collective effects are present:  $\Lambda = \frac{4\pi}{3}\lambda_D^3 n \gg 1$ ,
- electrostatic interactions dominate over  $2 \rightarrow 2$  scattering:  
 $\omega_p = \sqrt{\frac{4\pi\alpha n}{m}} \gg \Gamma_{2 \rightarrow 2}$ .

# Collisionless shocks

In counter-streaming plasma, electromagnetic instabilities can cause shock waves that lead to energy dissipation even if the mean free path determined by the  $2 \rightarrow 2$  scattering is much larger than the size of the system [1502.00626 [physics.plasm-ph]].



# Collisionless shocks

- Collisionless shocks are observed e.g. in the Earth's bow shock, in the expansion of supernova remnants into the interstellar medium and in X-ray emitting hydrogen plasma in galaxy collisions and cluster mergers.
- Collisionless shocks are studied numerically with particle in cell (PIC) simulations, and experimentally with electron-positron plasmas and ionized gases produced with laser pulses.
- Currently, numerical simulations of nonrelativistic pair-plasmas have not yet been performed.

# Contents

1 Introduction

2 Plasma dynamics

3 The model

# The model

- The goal is to explain the observed collisional behaviour of DM with energy dissipation caused by collisionless shocks.
- From observations of the bullet cluster, the fraction of collisional DM can be no more than 30%.
- In our minimal model we assume that 70% of DM is a generic WIMP, and 30% consists of *dark plasma*.

# Dark Plasma

- The minimal model for dark plasma is one Dirac fermion charged under an unbroken  $U(1)$  gauge group:

$$\mathcal{L} = \frac{1}{4} F_{D\mu\nu} F_D^{\mu\nu} + \bar{\chi} (i\not{D} - m_D) \chi.$$

- We neglect the kinetic mixing term  $F_{D\mu\nu} F^{\mu\nu}$  as it is highly constrained.
- The dark matter abundance is produced as a thermal relic by the annihilation into dark photons,  $\bar{\chi}\chi \rightarrow \gamma_D \gamma_D$ .

# Observational Constraints

- BBN bound on effective number of neutrino species  
 $N_{\text{eff}} < 3.38$  constrains the temperature of the dark photons during BBN.
- The dark photon temperature is given as

$$T_D = T_\gamma \left( \frac{g_{*s,\gamma}(T_\gamma) g_{*s,D}(T_*)}{g_{*s,D}(T_D) g_{*s,\gamma}(T_*)} \right)^{1/3},$$

where the two sectors are assumed to be in thermal equilibrium at  $T_*$ .

- This constrains the number of fermions in the dark sector:  
 $N_D < 2.35$ .



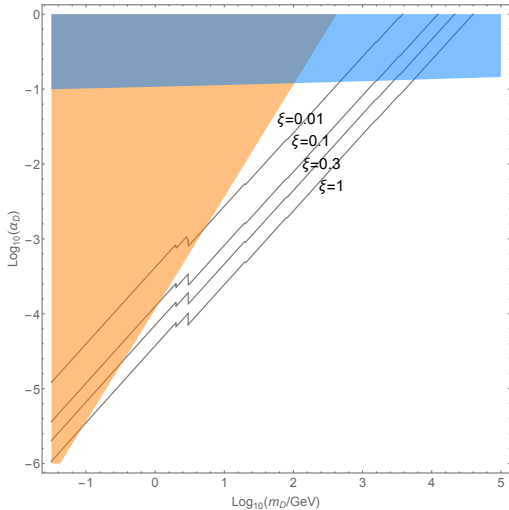
# Observational Constraints

- The rate of structure formation is suppressed until the kinetic decoupling of the dark matter and dark radiation, which occurs at

$$T_{\text{kin}} = \left( \frac{4\pi}{45} g_* \right)^{\frac{1}{4}} \sqrt{\frac{135}{64\pi^3}} \frac{m_D^{\frac{3}{2}}}{\sqrt{m_P \alpha_D}}.$$

- If  $T_{\text{kin}} > 640$  eV, only multipoles above  $l > 2500$  are affected in the CMB, and thus temperatures above this limit are unconstrained by PLANCK.
- For  $T_{\text{kin}} \approx 500$  eV the small scale structure is suppressed for structures below the size of  $\sim 10^9 M_\odot$ , alleviating the missing satellites problem.

# Results



# Outlook

- More observations required to form a coherent picture of DM dynamics in cluster mergers.
- More detailed analysis on dark plasma needed to understand its effects on galactic and cluster halos, structure formation etc.
- Further model building required for explaining naturally the partially interacting dark matter scenario, e.g. partially ionized dark atoms...